The Vainshtein mechanism in bimetric gravity arXiv:1506.04977

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Crhar Klein



Bimetric gravity

$$\begin{aligned} \mathcal{L} &= -\frac{M_g^2}{2}\sqrt{-\det g}R_g - \frac{M_f^2}{2}\sqrt{-\det f}R_f \\ &+ m^4\sqrt{-\det g}\sum_{n=0}^4\beta_n e_n\left(\sqrt{g^{-1}f}\right) + \sqrt{-\det g}\mathcal{L}_m \end{aligned}$$

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$$G_{\mu\nu}^{g} + m^{2} \sum_{n=0}^{3} (-1)^{n} \beta_{n} g_{\mu\lambda} Y_{(n)\nu}^{\lambda} \left(\sqrt{g^{-1}f}\right) = \frac{1}{M_{g}^{2}} T_{\mu\nu}$$
$$G_{\mu\nu}^{f} + m^{2} \sum_{n=0}^{3} (-1)^{n} \beta_{4-n} f_{\mu\lambda} Y_{(n)\nu}^{\lambda} \left(\sqrt{f^{-1}g}\right) = 0$$

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Graviton mass

$$m_g^2 = m^2(1+rac{1}{c^2})(eta_1 c + 2eta_2 c^2 + eta_3 c^3)$$

where $f_{\mu
u}=c^2g_{\mu
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- Runaway gravitational collapse in early universe?

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- Non-linear structure formation?

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Hubble radius

$$r_H \equiv H_0^{-1} \approx 5 \cdot 10^6 \, {
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$$r_V^3 \sim M \lambda_g^2$$

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Including higher order terms

$$\Phi = \Phi_{\mathrm{N}} \left[1 + k_1 \left(1 + \left(\frac{r_V}{r} \right)^3 + \cdots \right) \right]$$

• Expand in m_g^2

$$\Phi = \Phi_{\rm N} + k_2 \frac{r^2}{\lambda_g^2} + k_3 \frac{r^4}{\lambda_g^4} + \cdots$$
$$= \Phi_{\rm N} + k_2 \left(\frac{M}{r}\right) \left(\frac{r}{r_V}\right)^3 + k_3 \left(\frac{M}{r}\right)^2 \left(\frac{r}{r_V}\right)^6 + \cdots$$

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► To lowest order

$$\Phi = \Phi_{\rm N} \left[1 - k_2 \left(\frac{r}{r_V} \right)^3 \right]$$

• $r \lesssim r_V$

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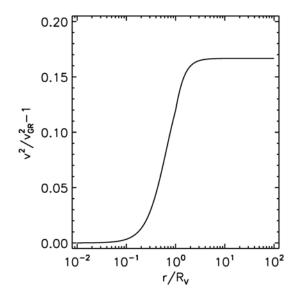
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• The Sun has $\Phi(r_V) = 10^{-15}$ and a galaxy $\Phi(r_V) = 10^{-7}$

Galaxy rotation curves



Deviations from GR in solar system small

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Galaxy lensing and velocity dispersions

$$\lambda_g \gtrsim 10^{-3} \, r_H$$
 or $\lambda_g \lesssim 10^{-6} \, r_H$

Vainshtein density

Average density within r_V of mass $M(r_V)$ is

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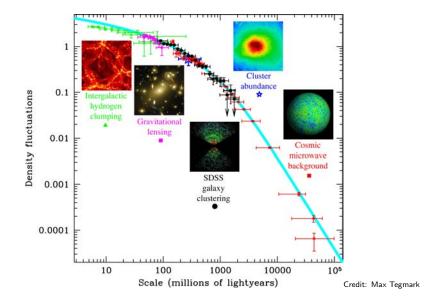
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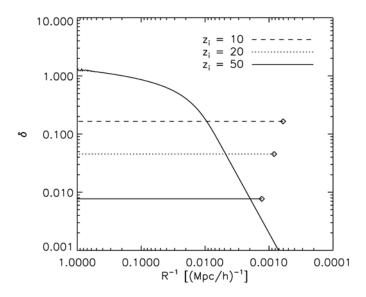
Fluctuation within its Vainshtein radius for

$$\delta \gtrsim \delta_V(a) \equiv \frac{\rho_V}{\bar{\rho}} \sim \frac{a^3}{\Omega_m}$$

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- Observations at large scales crucial